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25th

International Conference on Condensed Matter Nuclear Science

Developing a new energy source based on hydrogen-metal systems leading to clean, safe and low-cost high-density energy production, environmentally friendly and sustainable for economy and society.



SCIENTIFIC PROGRAMME:

- Heat Production
- Transmutations
- Electrochemical Experiments
- Hot gas experiments
- Plasma Experiments
- Beam Experiments
- Instrumentation
- Material Studies
- Theoretical and Computational Studies
- Engineering Applications
- Other

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$\sqrt[3]{A}$ —Law in Nuclear Transmutation of Metal Hydrides (II)



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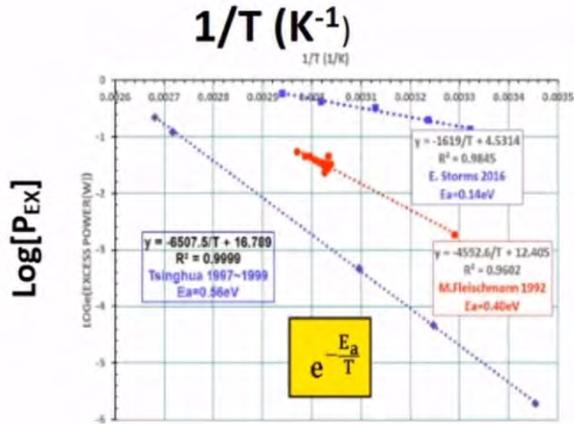
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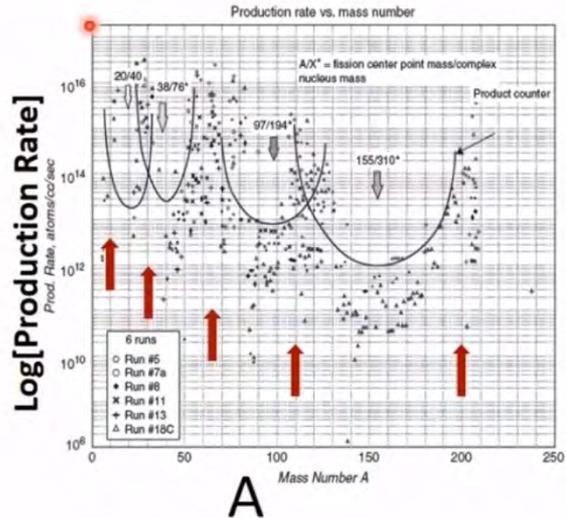
$\sqrt[3]{A_N}$ —Law in Nuclear Transmutation of Metal Hydrides (II)



$$\sqrt[3]{A_N} = 1.14 N - 0.39, \text{ Miley}$$

$$\sqrt[3]{A_N} = 1.13 N - 0.46, \text{ 2-Step, Resonance}$$

Elastic Facilitates Resonance, Resonance Enhances Reaction



CONCEPT: Resonance → Peaked Wave @ Nuclear Surface

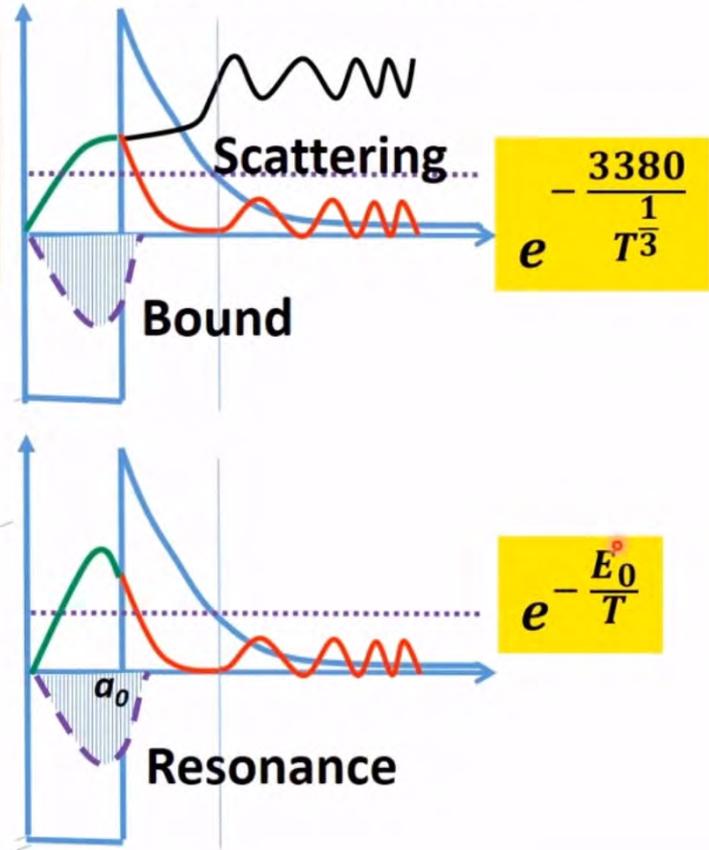


Solar Energy Model (Bethe 1938)

Storms: $\text{Log}[P_{\text{EX}}] = a \left(\frac{1}{T} \right) + b$

Smooth Connection:

$$a_0 \propto \left(\sqrt[3]{A} + 1 \right)$$



EQUATION: Logarithmic Derivative = That of Resonance Branch

EQUATION: Smooth Connection @ Nuclear Surface for Resonance

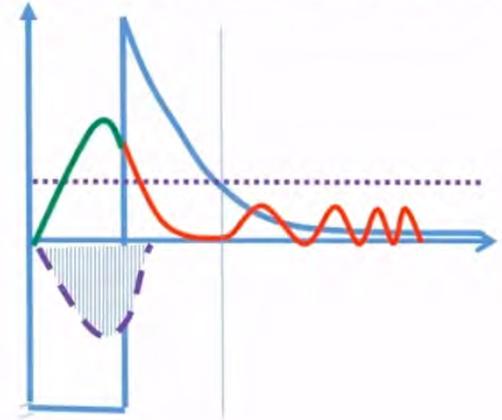
$$k \cot[k a_0] = - \sqrt{\frac{2}{a_0 a_c}}$$

$(\sqrt[3]{A} + 1) \quad \frac{1}{Z} \approx \frac{1}{A}$

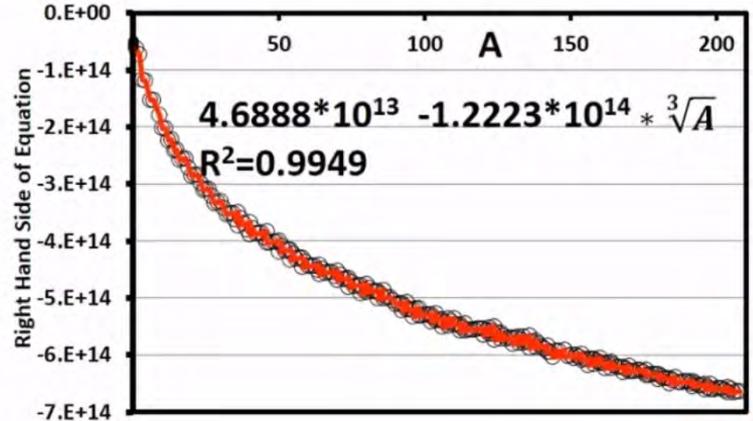
$$k [ka_0 - (2N + 1) \frac{\pi}{2}]$$

$$\downarrow$$

$$(\sqrt[3]{A} + 1) \quad -(1.22\sqrt[3]{A} - 0.469) \times 10^{14} m^{-1}$$



$$\sqrt[3]{A_N} = 1.13 N - 0.46, \quad \text{2-Step, Resonance}$$



COMPARISON: Isotope Depletion Peaks

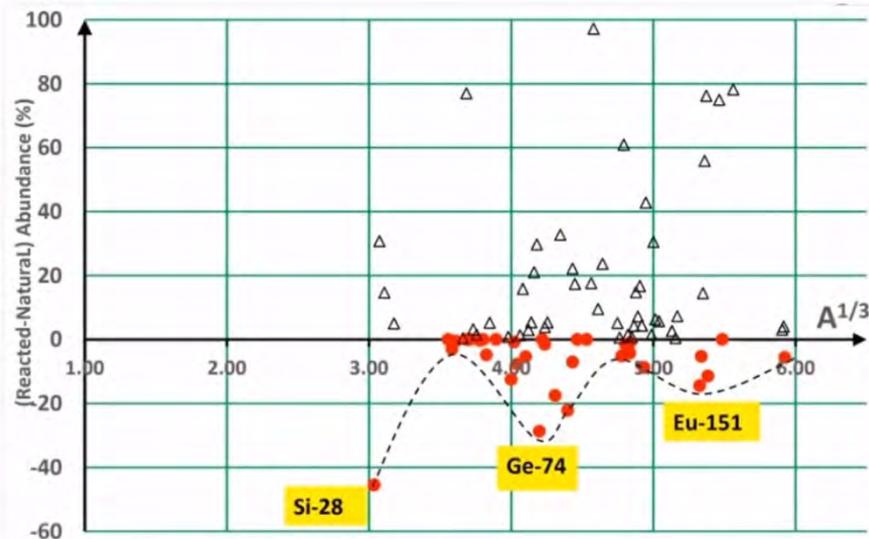
➤ $\sqrt[3]{A_N} = 1.14 N - 0.39$, Miley.

$\sqrt[3]{A_N} = 1.13 N - 0.46$,
2-Step Resonance.

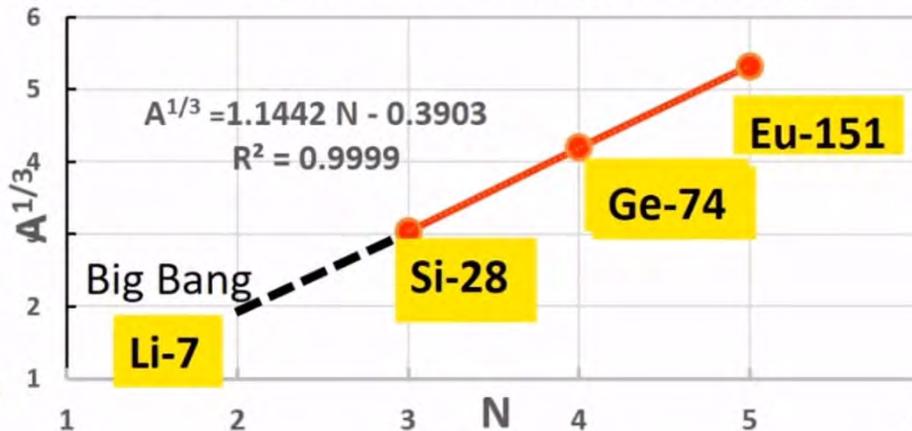
$a_0 \propto (\sqrt[3]{A} + 1)$

Target, Reactant, Not Product.

- 2 weeks. New Products were still in H. If an Isotope was Interacting with Proton. It was the Target, and would be depleted. Any Depletion Peak gives **Target A_N for Resonance**

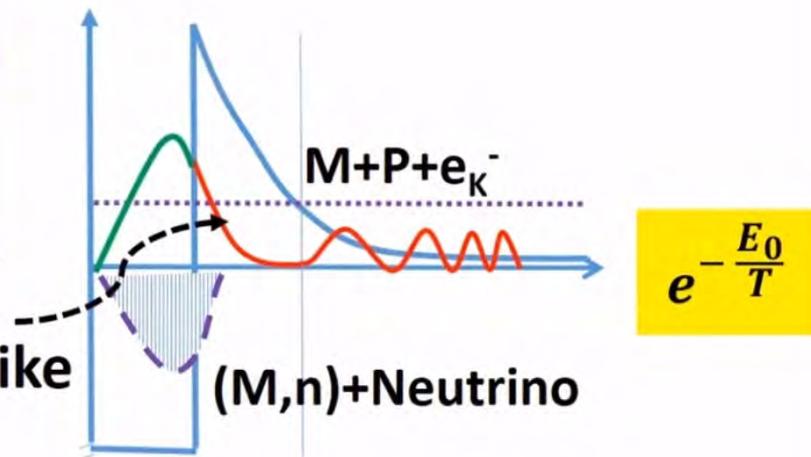
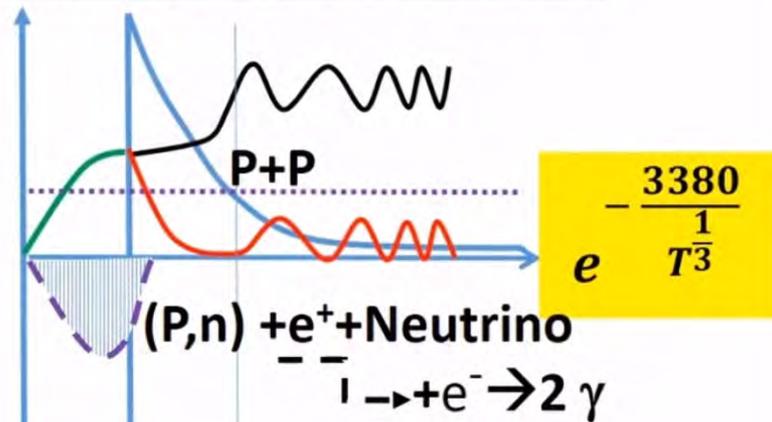


$A^{1/3}$ --Law in Nuclear Transmutation of Ni-H



NUCLEAR REACTION: K-Electron Capture vs. Positron Emission

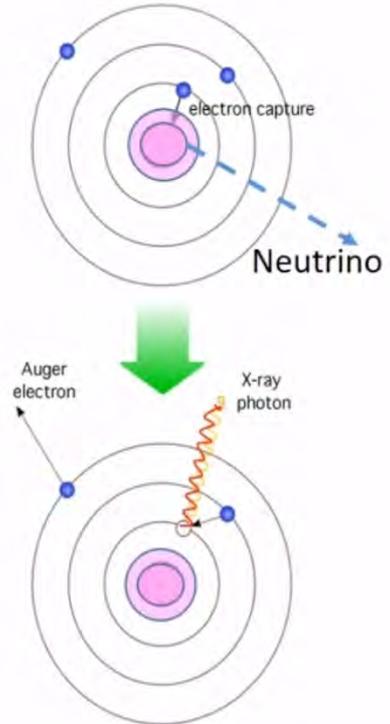
	Temperature	Charge	Energy Deficit (MeV)
SUN	~1.5keV, Gas ionized	Z~1	0.782 + 0.511
Miley	< 1 eV, K-Orbital Electrons	Z > 1, (Z+1) ³	0.782 - 0.511



Proton Halo-like

SUMMARY (M-H, M-D(1-rst Step))

- **Controlled Weak Interaction:**
No Neutron, No γ ;
Neutrino, Electrons, Recoil Ions, X-ray.
Gordon's Co-deposited Film (M-H).
- **2-Step Resonance:**
Miles(ICARUS), Letts, Mizuno, et al;
Miley(Ti-H 2003), Mizuno, Ohmori, Little&Puthoff.
- **Multiple-Scattering:**
Mckubre, Takahashi, Iwamura, Czerski, Kasagi,
Hagelstein, Meulenberg, Violante, Dardik, Vysotskii
- **Neutrino:**
 $1W \Rightarrow 1 M^2 \rightarrow 10^{15} / M^2 / Sec.$



Ardet nec consumitur

Потенциал Вудса — Саксона

Материал из Википедии — свободной энциклопедии

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Текущая версия страницы пока [не проверялась](#) опытными участниками и может значительно отличаться от [версии](#), проверенной 12 июля 2019 года; проверки требует [1 правка](#).

Потенци́ал Вудса — Саксона — функция [потенциальной энергии](#), предложенная^[1] американскими физиками [Д.Саксоном](#) и [Р.Вудсом](#), как приближение для той части потенциальной энергии [нуклона](#) в [атомном ядре](#), которая обусловлена ядерными силами и центрально-симметрична.

$$U(r) = -\frac{U_0}{1 + \exp\left(\frac{r-R}{a}\right)},$$

где

- $R = r_0 A^{1/3}$ — радиус ядра,
 - $r_0 \approx 1,25$ [фм](#), параметр, приближённо равный среднему расстоянию между нуклонами в ядре;
 - A — [массовое число](#) ядра,
- a — параметр диффузности, характеризующий размытие края [потенциальной ямы](#) (типичное значение — [0,5 фм](#)),
- U_0 — глубина потенциальной ямы (типичное значение — [50 МэВ](#)).



См. также [[править](#) | [править код](#)]

- [Оболочечная модель ядра](#)

Poster 21

- 5) [Первым найденным экзотическим решением является состояние с энергией $E_0 = -52.87$ эВ, орбитальный радиус $r_0 = 13.6 \cdot 10^{-12}$ м. При этом соответствующее решение уравнения Дирака имеет энергию $E_0 = -54.4$ эВ и орбитальный радиус $r_0 = 13.22 \cdot 10^{-12}$ м. Рис.2 [2].

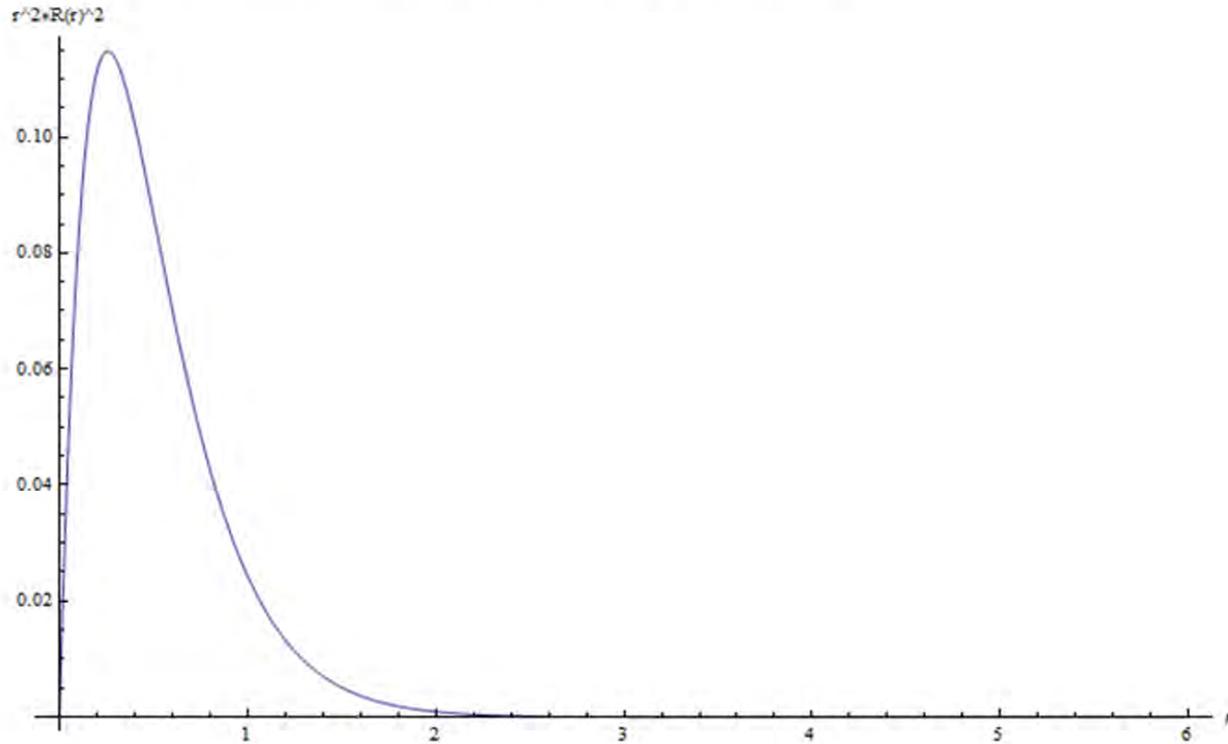


Рис.2 Нормированная радиальная плотность вероятности при $l = -\frac{1}{2}$ $E_0 = -52.8705 eV$

Последним найденным экзотическим состоянием является состояние с параметрами $E_0 = -mc^2 = -510785.13388932$ эВ и орбитальный радиус $r_0 = 0.006726620641 * 10^{-12}$ м. Рис.3. При образовании этого состояния выделяется энергия $E = mc^2 \approx 511$ кэВ. Выделение энергии $E = mc^2 \approx 511$ кэВ при образовании компактного атома водорода, может объяснить феномен появления аннигиляционного пика во многих экспериментах по НЭЯР [5] и в природных явлениях [6].

Косвенным доказательством существования компактного нейтроноподобного атома водорода, может служить эксперимент Дон Карло Борги [7].

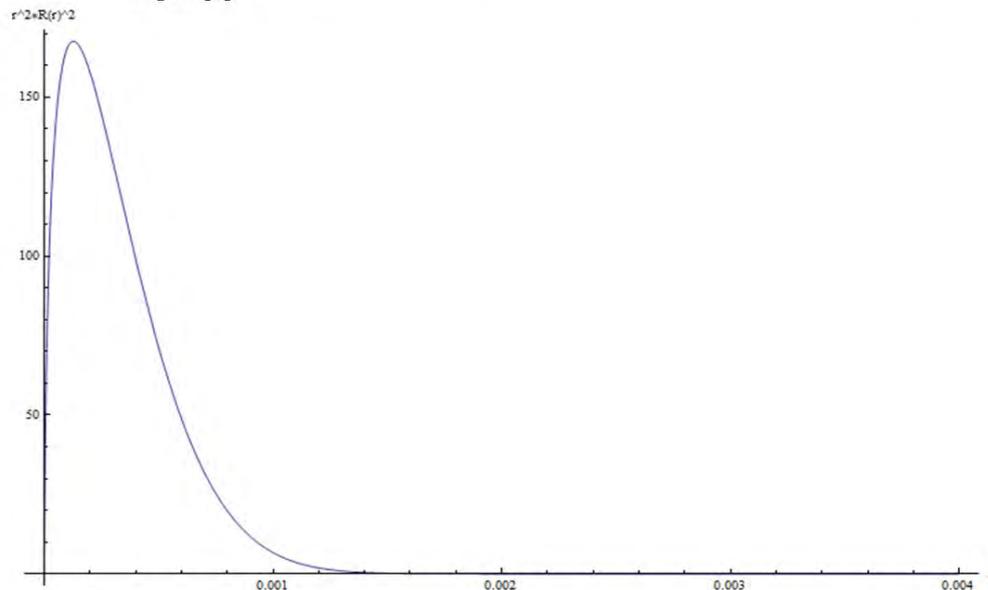


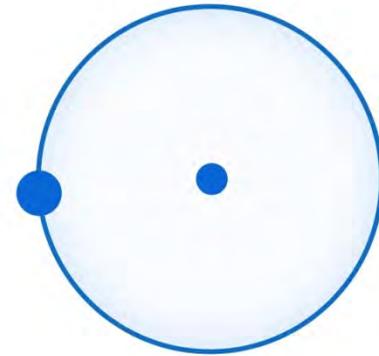
Рис. 3 Радиальная плотность вероятности при $l = -\frac{1}{2}$ и $E_0 = -510785.13388932$ эВ

Theoretical framework for LENR experiments.

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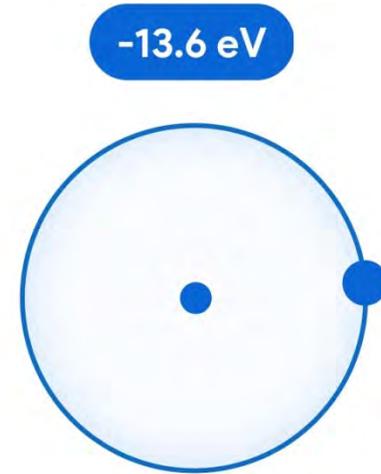
Hydrogen Atom

Theoretical framework for LENR experiments.

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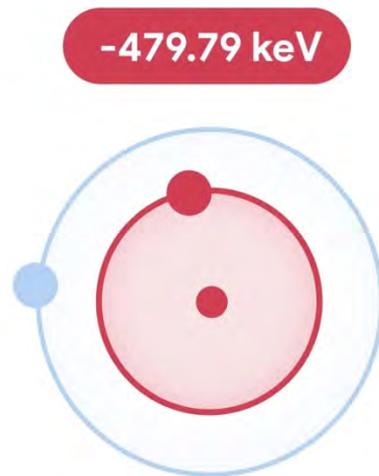
Ground State

Theoretical framework for LENR experiments.

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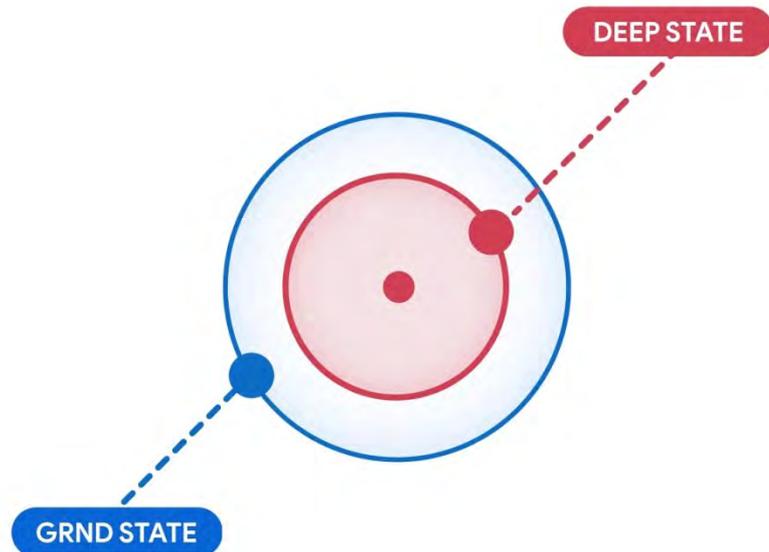
Deep State

Theoretical framework for LENR experiments.

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Exotic State

The radial M2 equation for a hydrogen atom.

$$\Delta\Psi - \frac{1}{\hbar^2} \left[\frac{m^4 c^6}{(E - U(\vec{r}))^2} - m^2 c^2 \right] \Psi = 0 \quad 1$$

Our journey begins with the relativistic M2 equation (*Equation 1*).

The radial M2 equation for a hydrogen atom.

$$\frac{d^2 R}{dr^2} + \frac{2}{r} \frac{dR}{dr} - \frac{l(l+1)}{r^2} R - \frac{1}{\hbar^2} \left[\frac{m^4 c^6}{\left(E + \frac{Ze^2}{4\pi\epsilon_0 r} \right)^2} - m^2 c^2 \right] R = 0$$

2

By integrating the Coulomb potential energy and separating variables, we derive the radial M2 equation, specifically tailored for hydrogen-like ions with nuclear charge **Z**.

The radial M2 equation for a hydrogen atom.

Next, we will apply the Hartree atomic unit system. In Hartree atomic units, physical constants have the following values:

$$a_0 = 1, m = 1, e = 1, \hbar = 1, c = 137.03599971, 4\pi\epsilon_0 = 1$$

Application of the Hartree Atomic Unit System

The radial M2 equation for a hydrogen atom.

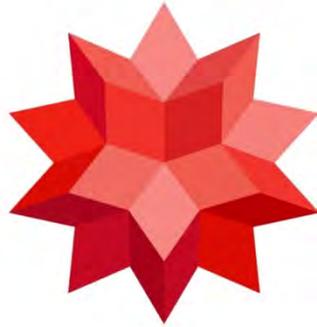
The radial M2 equation, now expressed in Hartree atomic units.

$$\frac{d^2 R}{dr^2} + \frac{2}{r} \frac{dR}{dr} - \frac{l(l+1)}{r^2} R - \left[\frac{c^6}{\left(E + \frac{Z}{r}\right)^2} - c^2 \right] R = 0$$

3

M2 equation in Hartree atomic units (*Equation 3*).

The radial M2 equation for a hydrogen atom.



Solution Approach using Wolfram Mathematica

The radial M2 equation for a hydrogen atom.

$$R(r) = \frac{1}{r} k_1 \exp\left(-\frac{r\sqrt{c^6 - E^2 c^2}}{E}\right) (Er+Z)^{\frac{1}{2}} \sqrt{4L + \frac{4Z^2 c^6}{E^4} + 1} {}_2F_1\left(\frac{1}{2}, \sqrt{4L + \frac{4Z^2 c^6}{E^4} + 1}, \frac{c^6 Z}{E^2 \sqrt{c^6 - E^2 c^2}}; \sqrt{4L + \frac{4Z^2 c^6}{E^4} + 1}, \frac{2\sqrt{c^6 - E^2 c^2} (Er+Z)}{E^2}\right)$$

4

The solution to the radial equation is a sum of two linearly independent parts. Generalized Laguerre polynomials and a degenerate hypergeometric function of the second kind.

We will use the second linearly independent solution.

The radial M2 equation for a hydrogen atom.

$$\frac{1}{2} + \frac{\sqrt{4l(l+1) + \frac{4Z^2 c^6}{E^4} + 1}}{2} - \frac{c^6 Z}{E^2 \sqrt{c^6 - E^2 c^2}} = -n_{rad}$$

$$E = \pm \frac{c \sqrt{(c + 2cn)^2 - Z^2}}{1 + 2n} \quad \text{6}$$

EXCITED STATE

$$E_0(Z) = c \sqrt{c^2 - Z^2} \quad \text{7}$$

GROUND STATE

The quantum number takes values from the series $n = 0, \frac{1}{2}, 1, \frac{3}{2}, 2, \frac{5}{2}$ with a step of $\frac{1}{2}$.

When the quantum number is $n = 0$, (**Formula 6**) transitions to the formula for the ground state energy equation (**Formula 7**).

The radial M2 equation for a hydrogen atom.

$$E = 27.2 \left(\frac{c \sqrt{(c + 2cn)^2 - Z^2}}{1 + 2n} - c^2 \right) \text{ eV}$$

8

$$E_0(Z) = 27.2 \left(c \sqrt{c^2 - Z^2} - c^2 \right) \text{ eV}$$

9

The energy obtained includes the electron rest energy mc^2 . Taking this into account, we will finally write the binding energy formulas in electron-volts. We will use only positive energy values, although the M2 equation gives symmetrical solutions.

The radial M2 equation for a hydrogen atom.

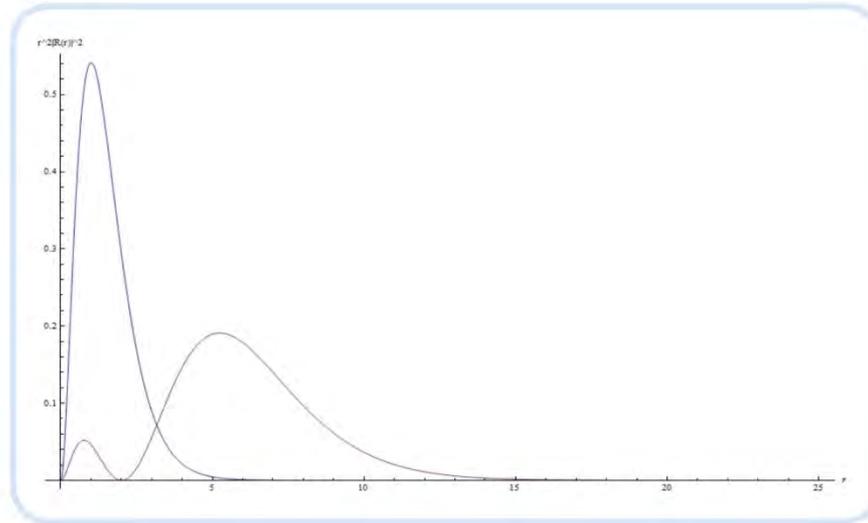


Figure 1. Normalized radial probability density of the ground and first excited states in Hartree atomic units.

Compact, deeply localized states of the hydrogen atom.

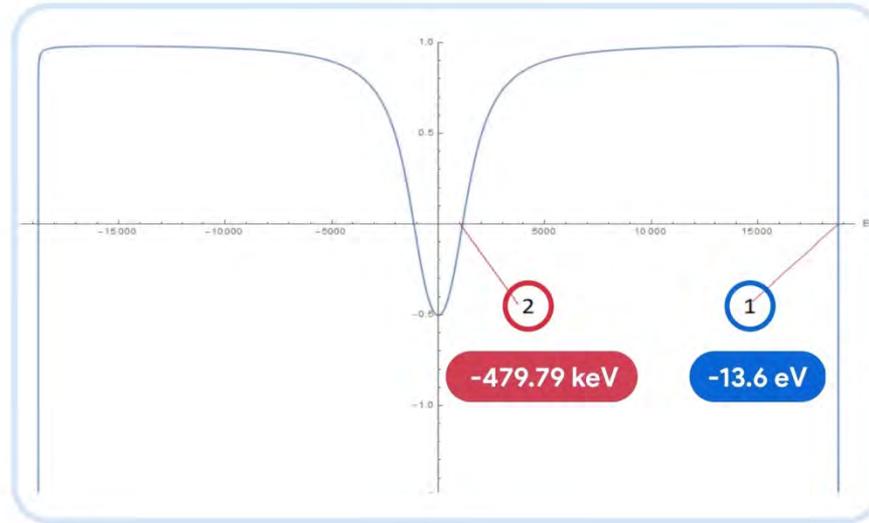


Figure 2. Graph of energy dependence on the radial quantum number $-n_{rad}$ at $l = 1$ for the M2 equation.

Compact, deeply localized states of the hydrogen atom.

$l \setminus n$	0	-1	-2	-3	-4	-5
0	-13.60					
1	-3.40	-479790.38				
2	-1.51	-492901.40	-472899.07			
3	-0.85	-498141.38	-486843.79	-470863.05		
4	-0.54	-500991.93	-492944.69	-484672.95	-469882.58	
5	-0.38	-502789.26	-496482.52	-490847.08	-483541.22	-469305.29

Table 1. Binding energies of compact hydrogen atom as a function of quantum numbers (in eV).

Radial Probability Density

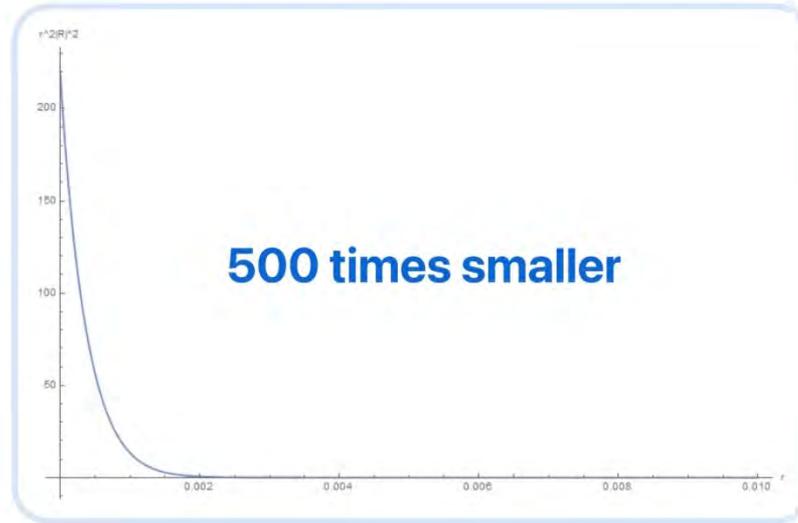


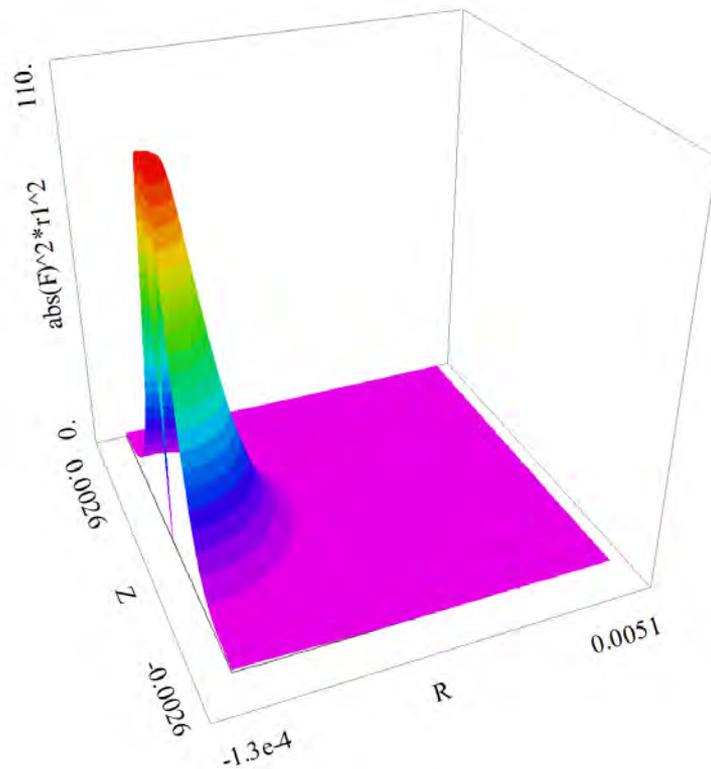
Figure 3: The radial probability density of state 2 demonstrates the potential for compact states in hydrogen atoms, a pivotal discovery in our study.

Решение уравнения M2
в цилиндрической системе координат

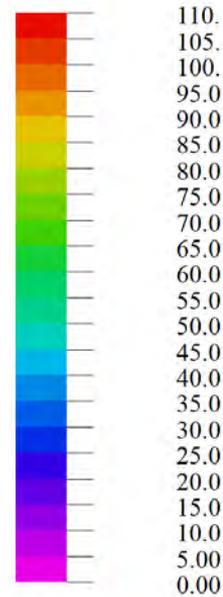
M2 Equation in Cylindrical Coordinates

$$\frac{\partial^2 R}{\partial r^2} + \frac{1}{r} \frac{\partial R}{\partial r} + \frac{\partial^2 R}{\partial z^2} - \left[\frac{c^6}{\left(E + \frac{Z}{\sqrt{r^2 + z^2}}\right)^2} - c^2 \right] R = 0 \quad 10$$

Equation 10: M2 equation in cylindrical coordinates in Hartree atomic units



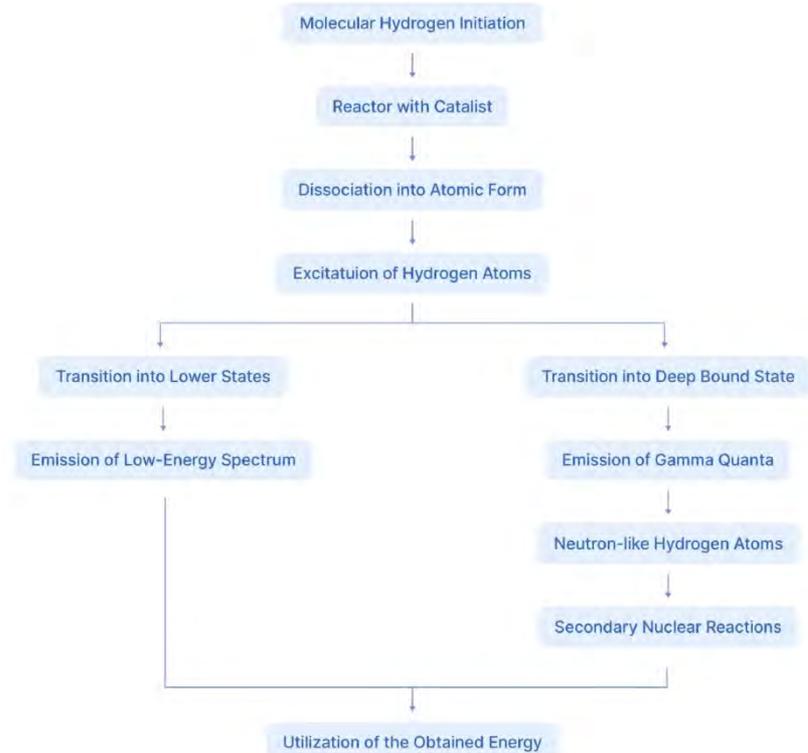
$\text{abs}(F)^2 * r1^2$
(-6.02e-3, -0.02, 30.)



M2Z=1: Grid#6 p2 Nodes=12501 Cells=6148 RMS Err= 2.e-7
E0 eV= -479790.4 Orbital Radius*52.9 (pm)= 0.010445 Vol_Integral= 4.246931e-7

Building on the deep state hydrogen atom theory:

1. Molecular hydrogen is introduced into a catalyst-filled reactor.
2. The hydrogen, upon interacting with the catalyst, transitions into its atomic form.
3. We then excite these hydrogen atoms.
4. During this excitation, some atoms transition into a deep bound state, emitting energy as gamma quanta.
5. Others move to lower states, releasing both gamma quanta and a typical low-energy spectrum.
6. These compact, neutron-like hydrogen atoms can participate in secondary nuclear reactions.
7. We then harness this released energy using various methodologies.



References:

- [1] Kurchatov, I. V. (1956). On the possibility of creating thermonuclear reactions in a gas discharge. Atomic Energy Journal, (3).
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- [5] Fu, C., Zhang, G., & Ma, Y. (n.d.). New opportunities for nuclear and atomic physics on the femto- to nanometer scale with ultra-high-intensity lasers. doi: 10.1063/5.0059405.
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- [7] 25 Конференция по Холодной Трансмутации Ядер и Шаровой Молнии 18 октября 2018. Доклад Дангыяна А.Э. "Экзотические состояния атома водорода"
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- [8] Poster21 ICCF 25 Danghyan Arayik <https://www.youtube.com/watch?v=GaGlp3o7Fp8>

THANK YOU!



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